# 7

## Carbon Redox Chemistry: Deep Carbon Dioxide and Carbonates

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#### ABSTRACT

Carbon dioxide is an important terrestrial volatile, often considered to exist in deep Earth interior, and it plays a critical role in the deep carbon cycle. Yet it is not well understood whether or how such volatile carbon species are incorporated into the Earth's lower mantle and core. Here, we present the phase/chemical diagram of carbon dioxide that signifies the transformation of molecular  $CO_2$  to extended solids, ionic carbonates, and ultimately its elemental decomposition products of C and  $O_2$ . We also discuss fundamental barochemical concepts governing carbon redox processes that may occur to recycle carbonaceous species in deep Earth's interior.

## 7.1. INTRODUCTION

Carbon dioxide exists in the atmosphere (as a major greenhouse gas), the ocean, and near the surface of Earth's crust, primarily as carbonates such as calcite and dolomite (Zhu & Ogasawara, 2002). It is also considered to exist deep in the Earth's mantle (Saal et al., 2002) and play a critical role in the deep carbon cycle (Hazen et al., 2013). The presence of carbon dioxide deep in Earth's interior is apparent from its abundance in magma plumes (Schrauder & Navon, 1993) and from assemblages of carbonate minerals such as aragonite and magnesite (Isshiki et al., 2004). In various thermal and chemical conditions, carbon dioxide converts to a wide variety of carbonaceous species, including diamond, graphite, carbon monoxide, and carbonates. Thus, the transformation of carbon, carbon dioxide, and carbonates at high pressures and temperatures is critical to understanding the origin of deep carbon, carbon budget, and carbon cycle (Hazen et al., 2013; Mao et al., 2003; Sleep et al., 2001).

Chemistries and properties of carbonaceous volatiles such as C, CO, CO<sub>2</sub>, and carbonates are not well understood at high pressures and temperatures (P-T). Yet, it is not well understood whether or how such volatile species are incorporated into the Earth's lower mantle and core,

resulting in a major uncertainty in the deep carbon cycle. The presence of CO, deep in the Earth can alter the properties of minerals and rocks, including the melting temperature and mechanical strength and, thereby, strongly affect the chemical and physical stability of deeply subducting plates (Dasgupta & Hirschmann, 2006; Kelemen & Manning, 2015). The mass and thermal diffusion rates of carbon dioxide into rocks and minerals are important considerations for carbon dioxide sequestration (Kampman et al., 2014; Leckner, 2003), but they are largely unknown at elevated P-T conditions, leaving questions about longterm storages in depleted oil and gas reservoirs. The relative stabilities of carbonate minerals (such as calcite CaCO<sub>3</sub>, magnesite MgCO<sub>3</sub>, and siderite FeCO<sub>3</sub>) with respect to the mixtures of high-density extended CO<sub>2</sub> (polymeric CO<sub>2</sub> or simply XCO<sub>2</sub>) and minerals (such as CaO, MgO, and Fe) are not well understood. Yet they are important in understanding both the deep mantle presence of volatile species like CO<sub>2</sub> and their geochemical transportation mechanisms. Importantly, the chemistry of CO, with other volatiles like  $H_2$ ,  $H_2O$ , or hydrous minerals is a primitive redox process, which can produce hydrocarbons and other simple organics and drive their evolutions.

Because of the significance of carbon dioxide in geochemistry, geological processes, and the deep carbon cycle, it is important to understand the phase and chemical behaviors of carbon dioxide at high P-T. Hence, we review

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*Carbon in Earth's Interior, Geophysical Monograph 249*, First Edition. Edited by Craig E. Manning, Jung-Fu Lin, and Wendy L. Mao. © 2020 American Geophysical Union. Published 2020 by John Wiley & Sons, Inc.

the phase/chemical diagram of carbon dioxide and present four governing, fundamental transformations: (A) carbon oxidation, (B) carbonate formation, (C) ionization of extended  $CO_2$ , and (D) pressure-induced amorphization. These results indicate a large P-T stability domain of extended  $CO_2$  and ionic carbonates, providing a geochemical mechanism for the origin of deep carbon species in the Earth's core and mantle boundary.

#### **7.2. EXPERIMENTAL METHOD**

The experiments were performed using laser-heated diamond anvil cells (DACs) coupled with confocal micro-Raman spectroscopy. Carbon dioxide samples were loaded into DACs from a liquid by condensing CO<sub>2</sub> gas to -35 °C and 15 atmospheres. Type IA diamond anvils were used with a flat culet size of 0.3 mm (or 0.1 mm) for higher pressures). A rhenium gasket was preindented to 0.04-0.05 mm in thickness, and a small hole of 0.12 (or 0.06) mm was drilled using an electric-discharge micro-drilling machine. A thin metal foil (Fe, Pt or B) (~0.01 mm) or a small piece of graphite was placed in the sample chamber to absorb the CW Nd: YLF laser ( $\lambda = 1054$  nm) light and heat the CO<sub>2</sub> sample. The temperature of the sample was determined by measuring thermal emission and then fitting it for a gray-body Planck radiation equation. The pressure of the sample was measured using the Ruby luminescence method (Mao et al., 1986). Existences and transitions of different carbon dioxide phases were determined based on their characteristic Raman spectra.

## 7.3. RESULTS

The phase diagram of carbon dioxide is rather intricate, with an array of polymorphs exhibiting great diversity in crystal structure, chemical bonding, and collective interaction (Figure 7.1) (Yoo et al., 2011). At pressures above 40 to 60 GPa and temperatures of 300 to 1000 K, carbon dioxide transforms into a wide range of silicate-like extended solids (or XCO<sub>2</sub>): four-fold CO<sub>2</sub>-V, pseudo-six-fold CO<sub>2</sub>-VI, coesite-like c-CO2, amorphous a-carbonia, and ionic carbonate solid *i*-CO<sub>2</sub>, all in high density between 3.5 and 3.8 g/cm<sup>3</sup>. These are fundamentally new solids, consisting of monolithic 3D network structures of carbon atoms covalently bonded with oxygen atoms, largely in CO<sub>4</sub> tetrahedra (or orthocarbonates) similar to those of silicate minerals. These solids exhibit novel properties such as extremely low compressibility (Datchi et al., 2012; Yoo et al., 1999), strong optical nonlinearity (Iota et al., 1999), and presumably high melting temperatures (as in the case of diamond). The large disparity in chemical bonding between the extended network and molecular CO<sub>2</sub>, on the other hand, allows these extended solids to exist over a large pressure domain (down to a few GPa), which provides opportunities for synthesis of novel materials.

To a first approximation, the crystallographic similarities between these extended  $CO_2$  phases and  $SiO_2$  polymorphs seem to indicate a periodic analogy between  $CO_2$ and  $SiO_2$  (Sengupta et al., 2010). Yet, unlike  $SiO_2$ , the polymerization of  $CO_2$  accompanies a great level of modification in chemical bonding from weak van der Waals in



**Figure 7.1** Phase and chemical transformation diagram of carbon dioxide, signifying the deep carbon cycle from molecular CO<sub>2</sub> phases to extended CO<sub>2</sub> solids (noted as XCO<sub>2</sub>), ionic solids (i-CO<sub>2</sub>) and, ultimately, decomposition products of C+O<sub>2</sub> at elevated temperatures and pressures. The blue line signifies the pressure (and temperature)– induced ionization of CO<sub>2</sub> phases, established by the specific transformations along the four blue arrows and described in sections 7.3.1 through 7.3.4. See electronic version for color representation of the figures in this book.

molecular solids ( $CO_2$ -I,  $CO_2$ -III, and  $CO_2$ -VII) to strong covalent bonds in extended  $CO_2$ . As a result, the phase diagram of carbon dioxide shows several thermalpath dependent phases (e.g., phase VI, c-coesite, and to an extent phase II and IV) and phase boundaries, underscoring the metastability and kinetics over a large P-T domain, making it difficult to determine the exact location of phase boundaries either experimentally or theoretically (Yoo, 2013).

Extended CO<sub>2</sub>-V exhibits a stability over a large P-T domain: to 2000 K at 60 GPa and to 200 GPa at ambient temperature, above which it transforms to extended ionic solids (*i*-CO<sub>2</sub> or COCO<sub>3</sub>) in the structure similar to postaragonite CaCO<sub>3</sub> (Ono et al., 2005) and, ultimately, decomposes to elemental solids of carbon and oxygen at extreme temperatures above 4000 K (according to the previous shock wave studies (Nellis et al. 1991). Recently, the stability of CO<sub>2</sub>-V has been reported even for a large P-T space well into the ionization line in Figure 7.1 (Dziubek et al., 2018), which may suggest a strong kinetic barrier associated with the transition. The presence of extended covalent (XCO<sub>2</sub>) and ionic (i-CO<sub>2</sub>) solids over a large P-T regime that overlaps with the Earth geotherm has significant implications to understanding the geochemistry of carbonaceous volatiles (C, CO, CO, and carbonates) and deep Earth's carbon cycle. Hence, described in sections 7.3.1 through 7.3.4 are governing carbon transformations along the specific thermal paths marked with arrows in Figure 7.1.

## 7.3.1. Carbon Oxidation

Carbon dioxide molecules are extremely stable at low pressures below 20 GPa, even to high temperatures above 3000 K. This is in contrast to the polymerization of CO<sub>2</sub>

above 40 GPa. It clearly indicates the presence of a strong kinetic barrier for breaking C=O bonds at low pressures than at high pressures. In fact, CO<sub>2</sub>-V can be made at a substantially lower pressure than 20 GPa, by heating phase IV in the presence of Ti catalyst (Sengupta et al., 2012). The polymerization of  $CO_2$  is, on the other hand, an oxidation process converting  $sp^2$  hybridized carbon in C=O to  $sp^3$  carbon in C-O. As such, it is thought that the presence of oxygen can assist carbon oxidation and the formation of new CO, phases, especially at high temperatures. Figure 7.2 shows the catalytic reaction between hot carbon particles and oxygen molecules, producing CO, and  $CO^{2+}CO_{3}^{2-}$ , an ionic form of  $(CO_{2})_{2}$  at 2300 K and 8 GPa. The spectral evidence includes the characteristic vibrational peaks of CO<sub>3</sub><sup>2-</sup> at 740 and 1075 cm<sup>-1</sup> and CO<sup>2+</sup> at 2230 cm<sup>-1</sup> for COCO<sub>3</sub>, and the Fermi resonance  $\nu$  and  $\nu_{\perp}$  peaks of CO<sub>2</sub> on either sides of the diamond Raman peak at ~1335 cm<sup>-1</sup>. Similar ionization reactions have previously been observed in nitrogen dioxide dimer  $(N_2O_4)$ both at high P-T (Agnew et al., 1983) and at ambient pressure and low temperature (Bolduan et al., 1984). It transforms to NO<sup>+</sup>NO<sub>2</sub>, whose vibrational spectrum is analogous to that of COCO<sub>3</sub>. Figure 7.3 shows a striking similarity in lattice vibrations between COCO<sub>2</sub> and  $NONO_2$  as an example.

These results clearly indicate the occurrence of carbon oxidation:  $C + O_2 \rightarrow CO_2 \rightarrow COCO_3$  at relatively low pressures and high temperatures, thus overcoming the kinetic barrier. Note that  $COCO_3$  is still a molecular ionic phase in contrast to extended ionic phase *i*-CO<sub>3</sub> formed at substantially higher pressure above 80 GPa (see Figure 7.1 and discussion in section 7.3.3). Nevertheless, the transformation occurs at the onset of the same phase line (the blue line in Figure 7.1), underscoring high stability of ionic solids over a large P-T region of the Earth mantle



**Figure 7.2** Raman spectra of laser-heated C and  $O_2$  mixtures to 2300 K at 10 GPa in the spectral regions for (a) internal vibrations and (b) external lattice modes. It shows the evidences for formation of  $CO_2$  and  $COCO_3$  (an ionic form of  $CO_2$  dimer). See electronic version for color representation of the figures in this book.

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**Figure 7.3** (a) Microphotograph of  $CO_2$ -Fe samples before (top, at ambient temperature) and during (bottom, at 2228 K) laser heating at 35 GPa in DAC. (b) Raman spectra of laser-heated  $CO_2$  through Fe to 2230 K at 35 GPa, showing the transformation of  $CO_2$ -III to  $CO_2$ -II,  $CO_2$ -V, and FeCO<sub>3</sub> as temperature increases. See electronic version for color representation of the figures in this book.

and core. Thus, this result highlights the significant role of extended carbonates in the deep carbon cycle.

#### 7.3.2. Carbonate Formation

The formation of carbonates at high temperatures and low pressures (Figure 7.2) advocates a potential geochemical mechanism of transporting carbonaceous volatiles into deep Earth interiors (Manning et al., 2013). In fact, carbonate minerals (especially of Ca and Mg) are abundant mostly in the Earth's crust (Oganov et al., 2013). To examine the formation of carbonate in deep Earth's mantle, we have investigated the transformation of CO<sub>2</sub> on the presence of Fe abundant in Earth's deeper mantle and core. Figure 7.3a shows the microphotograph of highly strained CO<sub>2</sub>-III and Fe samples in DAC, before (top at ambient temperature) and during (bottom at 2228 K) laser heating at 35 GPa. Raman spectra of CO<sub>2</sub> laser heated through Fe were taken at several temperatures at the same pressure, as noted on Figure 7.3b. The results indicate the transformation of CO<sub>2</sub>-III to initially CO<sub>2</sub>-II at 776 K and then CO<sub>2</sub>-V at 1823 K, as well as the formation of siderite FeCO<sub>3</sub> at 2228 K. The formation of FeCO<sub>2</sub> is apparent from the characteristic vibrational  $\nu_1$ peak of carbonate at ~1100 cm<sup>-1</sup> (Farsang, 2018). Note that the formation of siderite occurs at the onset of the ionization line of CO, in Figure 7.1, underscoring the significant role of siderite in the delivery of carbonaceous volatiles deep into the Earth's mantle.

## 7.3.3. Ionization of XCO<sub>2</sub>

The ionic character of the C-O bonds in XCO<sub>2</sub> increases further at high temperatures. Upon laser heating to ~1700-1800 K at 85 GPa, CO<sub>2</sub>-V transforms into an extended ionic solid (i-CO<sub>2</sub>), characterized by four characteristic Raman bands at 2000, 1200, 850, and 400 cm<sup>-1</sup>, as shown in Figure 7.4a. Note that the systematic of this Raman spectrum is remarkably similar to those of ionic carbonate COCO<sub>3</sub> and previously observed ionic nitrate NONO<sub>3</sub> (Agnew et al., 1983; Bolduan et al., 1984). However, there are some important differences. For example, both the  $\nu_1$  (CO<sub>3</sub>) at 980 cm<sup>-1</sup> and particularly the  $\nu_{s}$  (CO) at 1950 cm<sup>-1</sup> appear at substantially lower frequencies than those of ionic solids at 1070 cm<sup>-1</sup> and 2220 cm<sup>-1</sup>, respectively, at ~10 GPa. Nevertheless, extrapolating the pressure dependence of the  $\nu_1$  (NO<sub>3</sub>) mode to 85 GPa greatly reduces the difference to within 20-30 cm<sup>-1</sup>, which can easily account for the mass difference. The substantially greater difference in the  $\nu_{\rm c}(\rm CO)$  mode cannot be explained in terms of pressure dependence or mass difference, but it may indicate a local structural difference of carbonyl ions in *i*-CO<sub>2</sub> from COCO<sub>3</sub>.

A plausible model for the large softening of  $\nu_s$  (CO) mode is an incorporation of CO ions to adjacent carbonate layers. This would then lead to a structure analogous to the theoretically predicted mixed coordinated carbonate structures (Montoya et al., 2008; Sun et al., 2009). Hybridization of this type would certainly soften



**Figure 7.4** (a) Raman spectra of extended carbon dioxide phases before and after laser heating, showing the transformation to extended ionic  $CO_2$  carbonates (*i*- $CO_2$ ) at ~85 GPa and 1700–1800 K. (b) Raman spectra of extended carbon dioxide phases to 220 GPa, showing the pressure-induced amorphization of (top) four-fold  $CO_2$ -V via new high-pressure phase V' above 150 GPa, (middle) coesite- $CO_2$  (*c*-I) via a high pressure-form (*c*-II), and (bottom) phase III via *a*-carbonia, at ambient temperature. See electronic version for color representation of the figures in this book.

 $\nu_{s}$  (CO) and  $\nu_{1}$  (CO<sub>3</sub>) modes, as observed. In fact, the calculated vibration spectrum shows CO stretching in the range of 1800 to 1900 cm<sup>-1</sup>, depending on the degree of mixed coordination (Lee et al., 2009; Montoya et al., 2008). On the other hand, it is important to recognize that all calculated structures produce several strong extra bands between 700 and 400 cm<sup>-1</sup>, which are absent in *i*-CO<sub>2</sub> but present in disordered *a*-carbonia and phase VI. Therefore, the absence of such disordered peaks suggests a fully extended and more ordered *i*-CO<sub>2</sub> structure.

Note that *i*-CO<sub>2</sub> is formed only by heating extended CO<sub>2</sub> solids above 85 GPa; for example, heating phase III at 55 GPa produces only phase VIII (Sengupta & Yoo, 2009), whose Raman spectrum consists of two sharp peaks at 800 and 1200 cm<sup>-1</sup>, without the presence of the 2000 cm<sup>-1</sup> peak for CO.

#### 7.3.4. Pressure-Induced Amorphization

Above 40 GPa, carbon dioxide polymerizes to a wide range of covalently bonded extended solids: four-fold CO<sub>2</sub>-V (Datchi et al., 2012; Iota et al., 1999; Santoro et al., 2012; Sera et al., 1999; Yoo et al., 2013) and coesite-CO<sub>2</sub> (*c*-CO<sub>2</sub>) (Sengupta & Yoo, 2010), pseudo-six-fold CO<sub>2</sub>-VI (Iota et al., 2007), and mixed three- and four-fold coordinated *a*-carbonia (Santoro et al., 2006), each with a characteristic Raman-active  $\nu_{\rm b}$  (C-O-C) bending vibron at around 700–1000 cm<sup>-1</sup>. Upon further compression to 100–220 GPa, these XCO<sub>2</sub> phases become nonmetallic amorphous solids, as evident from complete loss of their vibrons and optical transparency. Figure 7.4b shows the corresponding spectral change of CO<sub>2</sub>-V to 220 GPa as an example.

Note that pressure-induced amorphization occurs at greatly diverse pressures, depending on extended phases: for example, CO<sub>2</sub>-VI at 80 GPa, *c*-CO<sub>2</sub> at 100 GPa, and CO<sub>2</sub>-V at 220 GPa. In fact, intermediate phases CO<sub>2</sub>-III and CO<sub>2</sub>-IV become amorphous solids at substantially lower pressures of 60 and 80 GPa, respectively (Santoro et al., 2006; Yoo et al., 2001). Yet it is remarkable that the pressure-induced amorphization occurs when the  $\nu_{\rm b}$  vibron reaches about the same Raman frequency, ~1000 cm<sup>-1</sup>, well above the  $\nu_{\rm b}$  (C-O-C) mode of six-fold configuration (CO<sub>6</sub>) at ~900 cm<sup>-1</sup>. Therefore, it is conceivable that

the amorphization represents a frustration of four-fold coordinated CO4 to increase either the packing density or the coordination number as pressure increases above 100 GPa.

A similar structural frustration or disorder was observed in CO<sub>2</sub>-VI, where carbon atoms are surrounded by an average of six oxygen atoms in a highly distorted octahedral with an average C-O distance of 1.45-1.71 Å (Iota et al., 2007), representing a substantial degree of ionic character in C-O bonds. Considering an approximately 10% in Si-O bond length between four-fold quartz (1.61 Å) and six-fold stishovite (1.76–1.81 Å) (Andrault et al., 1998), we speculate that the C-O bond length in CO<sub>2</sub>-VI must increase even larger to ~1.65–1.75 Å in order to accommodate six oxygen atoms around relatively small carbon atoms. With further compression, a separation of this size would eventually lead to a structural destruction forming amorphous solid, in which carbon atoms are in highly mixed coordination with six or more nearest neighbor atoms. The driving force is then to increase the packing density, as is apparent from the  $\nu_{\rm b}$ mode shifted to ~1050 cm<sup>-1</sup> (near that of carbonate). Therefore, it seems that the observed pressure-induced amorphization is driven by enhanced ionicity in carbonoxygen bonds and topological densification.

#### 7.4. DISCUSSION

The stability of CO<sub>2</sub> carbonates (both CO<sub>2</sub>-V and *i*-CO<sub>2</sub>) over a large pressure-temperature region relevant to the Earth's mantle and core strongly advocates a possibility that volatile carbon dioxide and carbonates are incorporated deep in the Earth's interior. Figure 7.5 supports this conjecture. It provides, for example, some constraints to evaluate the stability of carbonate minerals in descending slabs. It is also relevant to calculate the thermodynamics and phase assemblies associated with chemical reactions such as  $MgCO_3 \rightarrow MgO + CO_2$ and  $CaCO_3 \rightarrow CaO + CO_2$  (Dasgupta & Hirschmann, 2006). Using the equation of state of extended  $CO_2$ -V, for example, it can be readily shown that the mixture of extended CO, and MgO is substantially denser than MgCO<sub>3</sub> and other silicate minerals and is better matched to the Preliminary Reference Earth Model (PREM) in the upper mantle below 600 km (see Figure 7.5 inset). Therefore, even if MgCO<sub>3</sub> dissociates, CO<sub>2</sub> would remain either as extended CO<sub>2</sub>-V or *i*-CO<sub>2</sub> in descending slabs, delivering CO, deep into the Earth's mantle. A small contribution of iron in MgSiO<sub>3</sub> perovskite or postperovskite (Shim, 2008) can compensate for the density difference resulting from the presence of



**Figure 7.5** The pressure-density plots of various phases of XCO<sub>2</sub> and their analogs (C, NONO<sub>3</sub>), the major constituents of Earth's mantle (CaCO<sub>3</sub>, SiO<sub>2</sub>, MgCO<sub>3</sub>, and MgSiO<sub>3</sub>), and the Preliminary Reference Earth Model (PREM). Also plotted are the calculated densities of a ternary compound of two of the most common minerals in the Earth's core-mantle boundary (perovskite and wustite) and extended CO<sub>2</sub>:  $(1 - x)MgSiO_3 + xFeO + xCO_2-V$ , where x = 0.005, 0.01, and 0.02. The inset also shows the similar density variation of two binary compounds of extended CO<sub>2</sub> with MgO and CaO compared with two common minerals from the Earth's mantle, calcite and magnesite, and highlights the distinct possibility that MgCO<sub>3</sub> and CaCO<sub>3</sub> exist down to the 560 km transition zone of the Earth's mantle. See electronic version for color representation of the figures in this book.

extended CO<sub>2</sub>:  $(Mg_{1-x}, Fe_x)(Si_{1-x}, C_x)O_3 = (1 - x)MgSiO_3 + xCO_2 + xFeO, with <math>x = 0.5\%$  to 2.0% as it moves deeper into the core-mantle boundary from the 670 km transition zone. The decomposition of XCO<sub>2</sub> phases in the outer core region, then, provides a geochemical mechanism for the presence of deep carbon species (such as Fe<sub>3</sub>C) (Li et al., 1996; Nakajima et al., 2009) in the Earth's outer core, originally delivered from CO<sub>2</sub> in the atmosphere.

The P-T-induced transformations of CO<sub>2</sub> can be understood in terms of carbon redox processes that may have important implications to the deep carbon cycle (see Figure 7.6a). For example, the pressure-induced transformations of molecular  $CO_2$  to extended  $CO_2$ , covalent or ionic alike, are simply a carbon oxidation process, converting the  $sp^2$  C=O bonds in CO<sub>2</sub> to the  $sp^3$  C-O bonds in orthocarbonates  $(CO_4)$  or carbonates  $(CO_2)$  of XCO, phases. The pressure-induced ionization process, on the other hand, is a reduction process that converts the  $sp^3$  C-O bonds back to the  $sp^2$  C=O in *i*-CO<sub>2</sub> and further in C+O<sub>2</sub>. Such a reduction process is, then, assisted by high temperature. The rapid temperature rise at the Earth's mantle-core boundary can further accelerate the ionization and decomposition of extended CO<sub>2</sub> or carbonate minerals to form deep carbons and carbonaceous minerals. As such, it gives a rebirth of fully reduced carbon species deep in the Earth's mantle-outer core boundary region, where carbon oxidation can start again as it floats up over a geological time frame.

The pressure-induced transformation in dense  $CO_2$  also highlights several fundamental barochemistry (or pressureinduced solid-state chemistry) concepts, as illustrated in a hypothetical 2D lattice in Figure 7.6b (Yoo, 2017). At low pressures, molecular solids are relatively soft, primarily held by van der Waals interactions. Upon compression,

molecular solids transform into extended polymeric structures to soften highly repulsive intermolecular interaction by delocalizing valence electrons or making bonds between nearby molecules (i.e., the pressure-induced electron delocalization). These high-symmetry polymeric structures can further undergo symmetry-breaking distortions (such as Jahn-Teller or Peierls) to enhance their packing efficiencies, which result in low-symmetry distorted or even amorphous structures with highly polarized chemical bonds. Upon further compression, the electrostatic packing energy dominates over the electron hybridization energy, converting these covalent network structures into ionic solids (i.e., pressure-induced ionization) or extended mixtures of constituting elements (phase separation), depending on the chemical/bond/structural miscibility of dense solid interfaces of multiphase and multicomponents. These transformations are strongly controlled by kinetics, giving rise to the path dependence, metastable intermediate states, and near ground-state final products, as observed in CO<sub>2</sub> and many other molecular systems. The collective behaviors of lattice phonons, electrons, and charges in these extended structures, on the other hand, can give rise to novel electro-optical properties such as second harmonic generation (Iota et al., 1999) and high Tc superconductivity (Drozdov et al., 2015).

#### 7.5. CONCLUSION

We have described the transformations of carbon dioxide to extended solids, both covalent and ionic, at high pressures and temperatures. These results suggest the presence of volatile carbon dioxide and carbonates over a wide range of high P-T conditions in deep Earth's mantle, thus providing a geochemical mechanism explaining the origin of deep carbon species in the Earth's core



**Figure 7.6** (a) A schematic of pressure-induced transformation of solid  $CO_2$  in two dimensions, showing the fundamental concepts of pressure-induced chemistry (or barochemistry), such as the pressure-induced electron hybridization to extended solids and the pressure-induced ionization to ionic solids and decomposition products. (b) A schematic of deep carbon cycle governed by carbon Redox processes at combined high pressures and temperatures. See electronic version for color representation of the figures in this book.

and mantle boundary. In a context of carbon redox chemistry, we have conjectured a deep carbon cycle that may give a rebirth of fully reduced carbon species deep in the Earth's mantle–outer core boundary. We have also discussed fundamental barochemical concepts to explain the pressure-induced transformations in  $CO_2$  and other molecular solids.

#### ACKNOWLEDGMENTS

This work has been performed in support of the National Science Foundation, Division of Materials Research (Grant No. 1701360), the Department of Energy, National Nuclear Security Administration (DE-NA0003342), the Army Research Office (W911NF-17-1-0468), the donors of the Petroleum Research Fund administered by the American Chemical Society (No. 54806-ND10), and the Sloan Foundation through the Deep Carbon Observatory—Extreme Physics and Chemistry.

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